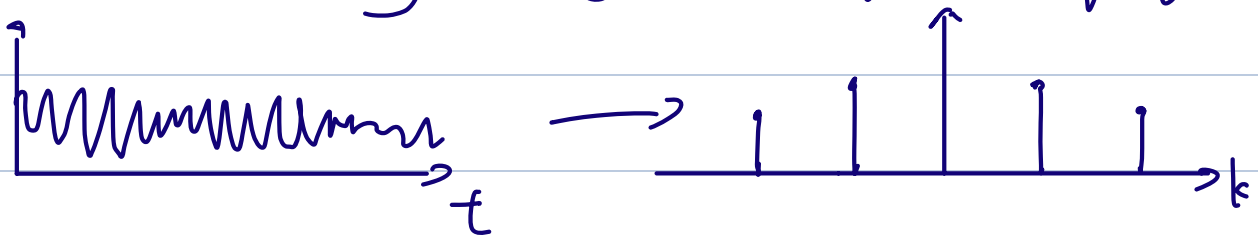


Previously, we showed that any periodic function can be represented as a series of harmonics:

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

where:  $a_k = \frac{1}{T} \int_0^T x(t) e^{-jk\omega_0 t} dt$

↑ Representation of a time dependent signal w/ harmonic sinusoids, where, the harmonic sinusoids are essentially integers at different frequencies.



But then, for the representation of the signal on "true frequency"  $\omega$ , &  $\because \omega = \frac{2\pi}{T} \cdot k = \omega_0 \cdot k$ , we must "normalize" amplitudes by  $T$  so that the total "energy" of the signal is consistent

Also Previously, the "harmonic condition" is matched for Amplitudes only when we use this assumption:

$$@ k=n, \int_0^T e^{j(k-n)\omega_0 t} dt = T, @ k \neq n, 0$$

When there is not a defined "T" in the original function, the  $\int_0^T e^{jk\omega_0 t} dt$  is not discrete at  $k=n$ : also  $k$  is continuous.

$\therefore$

$$\text{rewrite: } a_k = \frac{1}{T} \int_0^T x(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_{-\infty}^{\infty} x(t) e^{-jk\omega_0 t} dt$$

$$\text{define: } X(j\omega) = \int_{-\infty}^{+\infty} x(t) e^{-j\omega t} dt$$

$$\text{compare w/ } a_k: a_k = \frac{1}{T} X(jk\omega_0)$$

$$\therefore x(t) = \sum_{-\infty}^{\infty} \frac{1}{T} X(jk\omega_0) e^{jk\omega_0 t}$$

$$(T \rightarrow \infty) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega$$

$$\hookrightarrow X(j\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt$$

↑↑ Anyway, above is extra stuff for  
fourier transform, not in the scope of the  
course

Previously, in our string example:

$$y_n(x, t) = A_n \sin\left(\frac{n\pi}{L} x\right) \cos(\omega_n t)$$

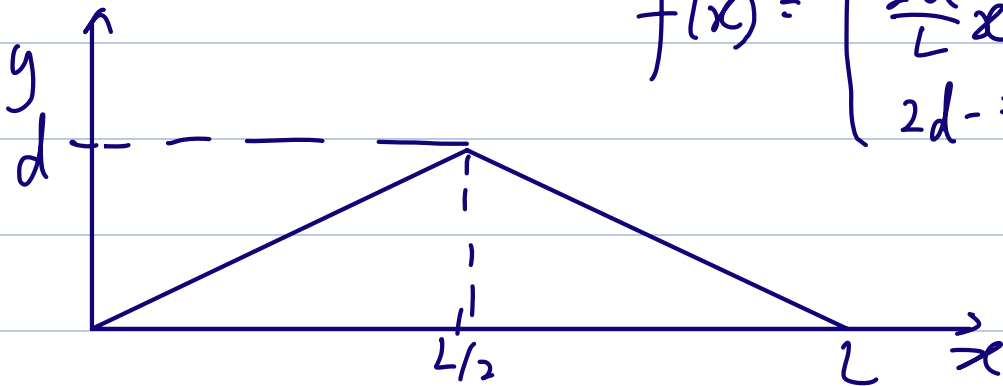
$$y(x, 0) = \sum_n A_n \sin\left(\frac{n\pi}{L} x\right) = f(x) \quad [\text{superposition of normal modes}]$$

Any string w/ fixed end points  $[f(0) = f(L) = 0]$  is:

$$\begin{cases} f(x) = \sum_n A_n \sin\left(\frac{n\pi}{L}x\right) \\ A_n = \frac{2}{L} \int_0^L dx \sin\left(\frac{n\pi}{L}x\right) f(x), n \in \{+\mathbb{Z}\} \end{cases}$$

example:

$t=0$ :



$$f(x) = \begin{cases} \frac{2d}{L}x, & x \in [0, L/2] \\ 2d - \frac{2d}{L}x, & x \in [L/2, L] \end{cases}$$

$$A_n = \frac{2}{L} \left[ \int_0^{L/2} dx \frac{2d}{L}x \sin\left(\frac{n\pi}{L}x\right) + \int_{L/2}^L dx \left(2d - \frac{2d}{L}x\right) \sin\left(\frac{n\pi}{L}x\right) \right]$$

$$= \frac{2}{L} \left[ \frac{2d}{L} \int_0^{L/2} x \sin\left(\frac{n\pi}{L}x\right) dx + 2d \int_{L/2}^L \sin\left(\frac{n\pi}{L}x\right) dx \right]$$

$$= \frac{2d}{L} \int_{L/2}^L x \sin\left(\frac{n\pi}{L}x\right) dx]$$

$$\int x \sin(ax) dx = \int u dv = uv - \int u du$$

$$u = x, du = dx$$

$$dv = \sin(ax) dx, v = -\frac{\cos(ax)}{a}$$

$$= -\frac{x}{a} \cos(ax) + \int \frac{\cos(ax)}{a} dx$$

$$= -\frac{x}{a} \cos(ax) + \frac{1}{a^2} \sin(ax)$$

$$\begin{aligned}
 \text{So: } \int_0^{L/2} x \sin\left(\frac{n\pi}{L}x\right) dx &= -\frac{L}{n\pi} x \cos\left(\frac{n\pi}{L}x\right) + \frac{L^2}{(n\pi)^2} \sin\left(\frac{n\pi}{L}x\right) \Big|_0^{L/2} \\
 &= -\frac{L}{n\pi} \frac{L}{2} \cos\left(\frac{n\pi}{2}\right) + \frac{L^2}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right) \\
 &\quad + \frac{L}{n\pi} 0 \cos(0) - \frac{L^2}{(n\pi)^2} \sin 0 \\
 &= \frac{L^2}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right) - \frac{L^2}{2n\pi} \cos\left(\frac{n\pi}{2}\right)
 \end{aligned}$$

$$\begin{aligned}
 \int_{L/2}^L x \sin\left(\frac{n\pi}{L}x\right) dx &= -\frac{L}{n\pi} x \cos\left(\frac{n\pi}{L}x\right) + \frac{L^2}{(n\pi)^2} \sin\left(\frac{n\pi}{L}x\right) \Big|_{L/2}^L \\
 &= -\frac{L}{n\pi} L \cos(n\pi) + \frac{L^2}{(n\pi)^2} \sin(n\pi) \\
 &\quad + \frac{L}{n\pi} \frac{L}{2} \cos\left(\frac{n\pi}{2}\right) - \frac{L^2}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right)
 \end{aligned}$$

$$\begin{aligned}
 \int_{L/2}^L \sin\left(\frac{n\pi}{L}x\right) dx &= -\frac{L}{n\pi} \cos\left(\frac{n\pi}{L}x\right) \Big|_{L/2}^L \\
 &= -\frac{L}{n\pi} \cos(n\pi) + \frac{L}{n\pi} \cos\left(\frac{n\pi}{2}\right)
 \end{aligned}$$

$$\begin{aligned}
 \text{plug in: } A_n &= \frac{2}{L} \left[ \frac{2d}{L} \cdot \frac{L^2}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right) + 2d \cdot \left(-\frac{L}{n\pi} \cos(n\pi)\right) \right. \\
 &\quad \left. - \frac{2d}{L} \left(-\frac{L^2}{n\pi} \cos\left(\frac{n\pi}{2}\right) + \frac{L^2}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right)\right) \right] \\
 &= \frac{2}{L} \left[ \frac{2dL}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right) + \frac{2dL}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right) \right]
 \end{aligned}$$

$$A_n = \frac{8d}{(n\pi)^2} \sin\left(\frac{n\pi}{2}\right)$$

Evaluation:

$$n = 1$$

$$A_n = \frac{8d}{\pi^2}$$

$$\sin\left(\frac{n\pi}{2}x\right)$$



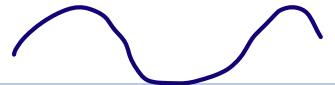
$$2$$

$$0$$



$$3$$

$$-8d/(3\pi)^2$$



$$4$$

$$0$$

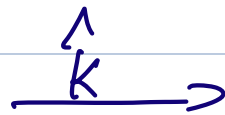
$$;$$

$$5$$

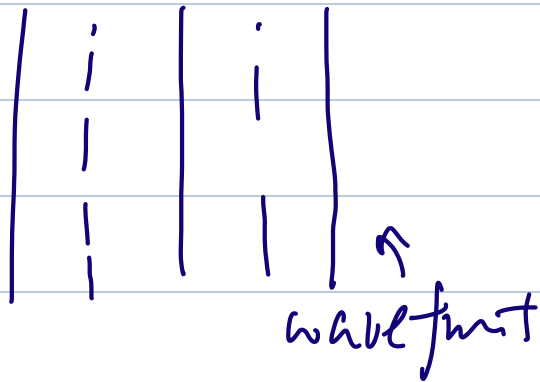
$$8d/(5\pi)^2$$

$$;$$

# Interference & diffraction:



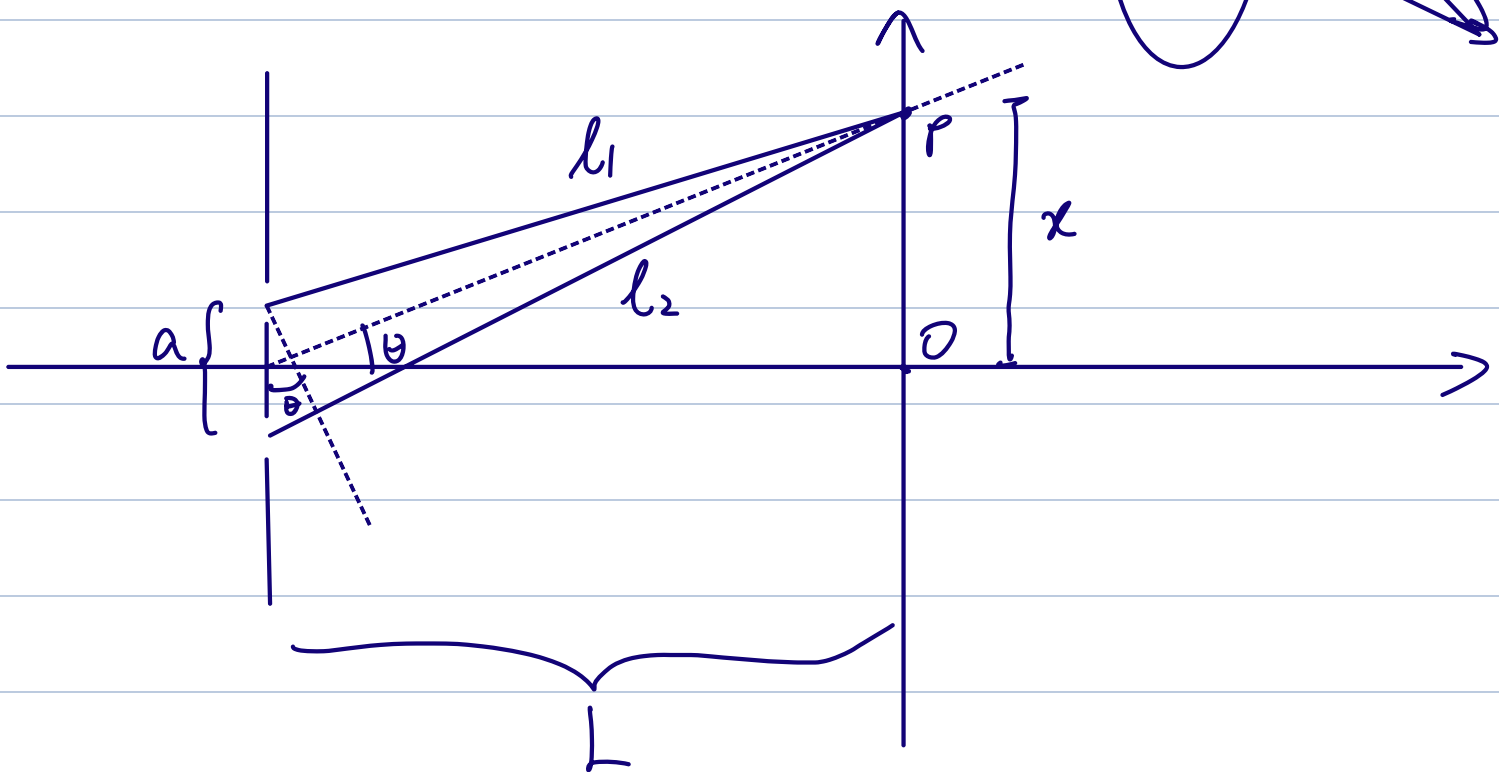
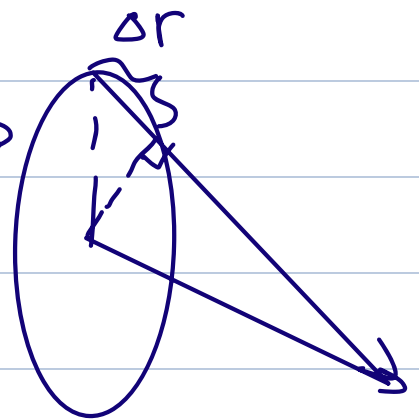
plane waves



Huygens:

any point on the wave front can be a source of a new wave

Diffraction is interference  $\rightarrow$   
at a small but dense scale



$$R = A [\cos(\omega t - k l_1) + \cos(\omega t - k l_2)]$$

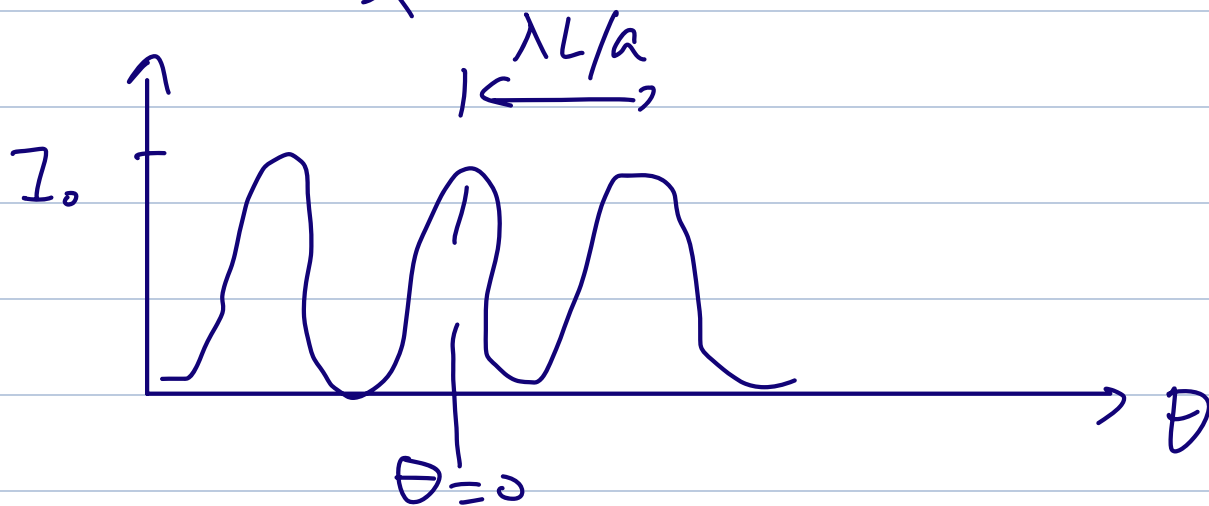
$$= 2A \cos(\omega t - k(l_1 + l_2)/2) \cdot \cos(k(l_2 - l_1)/2)$$

$$I \propto \frac{1}{2} R^2 = 2A^2 \cos^2[k(l_2 - l_1)/2]$$

$$l_2 - l_1 \triangleq \Delta l = a \sin \theta = a \theta$$

$$\therefore I(\theta) = I_0 \cos^2\left(\frac{k a \theta}{2}\right) = I_0 \cos^2\left(\frac{\pi a \theta}{\lambda}\right)$$

$$k = \frac{2\pi}{\lambda}$$



$$x = L \cdot \tan \theta \approx L \theta$$